

Advanced Transport Phenomena II
 ChE 548
 Department of Chemical Engineering
 University of Tennessee, Knoxville
 Spring, 2002

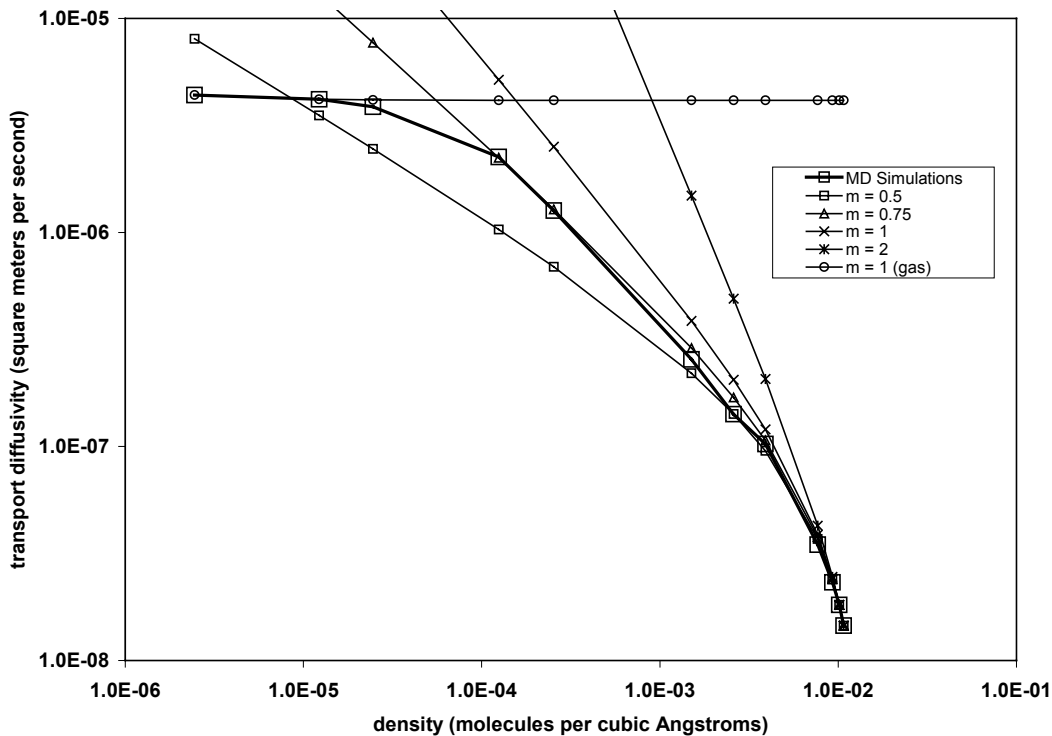
Final Examination
 Administered: May 2, 2002

Problem 1. (30 points)

Kinetic theory tells us that the diffusivity of an ideal gas is proportional to the inverse of the total molar density,

$$D = b_0 + b_1 \left(\frac{1}{n} \right)^m$$

so that in this case, the exponent, m , is one. (b_0 and b_1 are empirical constants.) Consider an equimolar binary mixture of methane and ethane at 300 K. The single, binary transport diffusivity, obtained by using MD simulations for the self-diffusivities and the Lennard-Jones equation of state for the thermodynamic correction factors, is plotted below as a function of total density (on a log-log scale). Also plotted are several asymptotes with the function form of the equation given above. At the low density (ideal gas limit), the $m = 1$ equation fits well. At the high density, several different values of m are plotted.



Answer the following questions based on this plot and on your understanding of the molecular-level mechanisms for diffusion in gases and liquids.

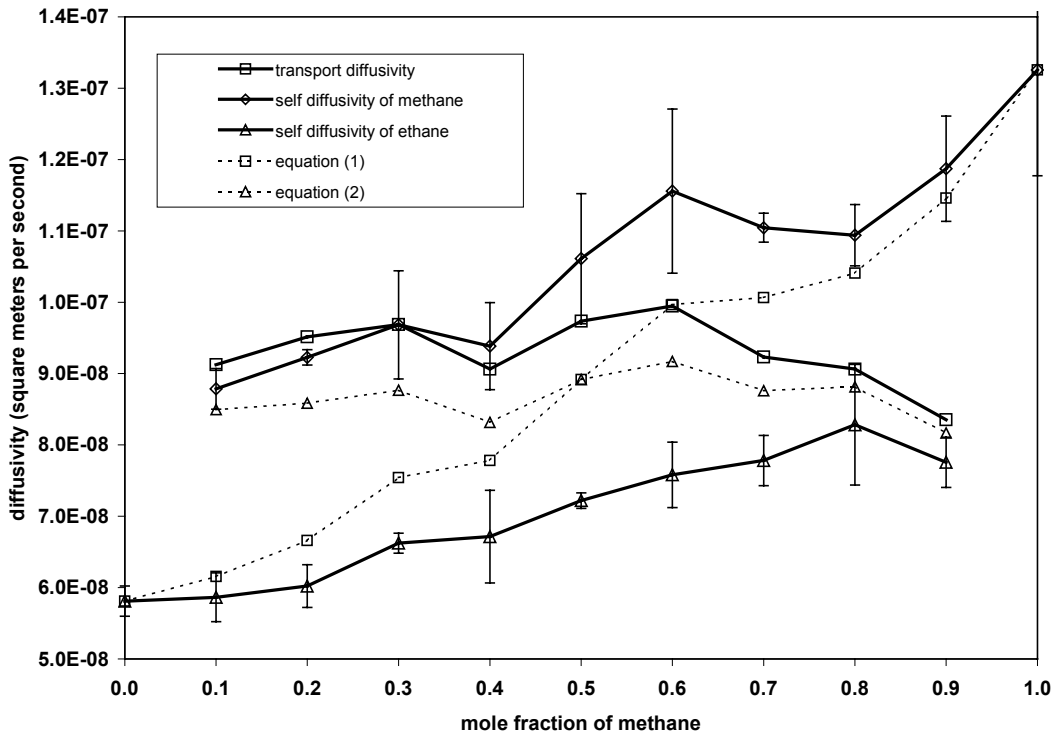
- What is the molecular-level mechanism for diffusion in gases that gives rise to $m=1$ behavior?
- What value of m models the high-density (liquid) density-dependence of the diffusivity best?
- What is the molecular-level mechanism for diffusion in liquids that gives rise to the m value selected in part (b)?

Problem 2. (30 points)

The single, binary transport diffusivity of a liquid is obviously a function of composition. Consider a binary mixture of methane and ethane at 300 K and at a density of 3.907×10^{-3} molecules/ \AA^3 (a liquid density). The single, binary transport diffusivity, obtained by using MD simulations for the self-diffusivities and the Lennard-Jones equation of state for the thermodynamic correction factors, is plotted below as a function of mole fraction of methane (on a log-log scale). Also plotted are the self diffusivities from MD simulations. Also plotted are two equations.

$$D = x_{\text{Me}} D_{\text{self,Me}} + x_{\text{Et}} D_{\text{self,Et}} \quad (1)$$

$$D = x_{\text{Me}} D_{\text{self,Et}} + x_{\text{Et}} D_{\text{self,Me}} \quad (2)$$



- Explain why the self diffusivity of methane is greater than the self diffusivity of ethane at the same temperature and density.
- The self diffusivities of methane and ethane generally increase with mole fraction of methane. Why?
- The self diffusivities of methane and ethane are not strictly monotonically increasing with mole fraction of methane. In your opinion, do the local non-monotonicities reflect real trends or, rather, are they within the noise of the simulation?
- In this data, the transport diffusivity is bound by the two self-diffusivities at all compositions. Must this behavior always be true?
- If you believe that the approximate physics of your simulations (such as the approximation of the Lennard-Jones potential) are only good to generate transport properties to $\pm 10\%$, is it a legitimate approximation to assume a constant mean value of the transport diffusivity for use in a process modeling application?
- The transport diffusivity is clearly better modeled by equation (2) than by the molar-weighted mean self diffusivity (equation (1)). Explain.

Problem 3. (40 points)

Consider a two-component ideal mixture with constant molar volume in a perfectly insulated one-dimensional system of length L . The boundary conditions of the system are such that at $z=0$, the composition is fixed at $x_{\alpha,0}$, and at $z=L$ the composition is fixed at $x_{\alpha,L}$. Since the system is insulated, there is no heat flux at the boundaries, $q = 0$. Initially the temperature is constant at the thermodynamic reference temperature. The enthalpy of component α at the reference temperature is positive (e.g. CH_4). The enthalpy of component β at the reference temperature is zero (e.g. N_2). In this system, no reaction takes place. You cannot ignore the enthalpy transfer due to the mass flux. (You must use the following form of the heat flux.)

$$q = -k \frac{\partial T}{\partial z} + \sum_{\alpha=1}^{n_c} \bar{H}_{\alpha} J_{\alpha}^*$$

As usual, q is the heat flux, k is the thermal conductivity, n_c is the number of components (2), \bar{H}_{α} is the partial molar enthalpy of component α , and J_{α}^* is the molar flux of component α relative to the molar average velocity, v^* .

In this problem, you may assume that the thermal conductivity and diffusivity are constants with respect to temperature, density, and composition. Assume the convection term is zero at steady state.

- Express q for an ideal mixture.
- Write out the total unsteady state mole balance, the balance on moles of α , and the energy balance.
- Derive the steady state profile of the concentration.
- Derive the steady state profile of the mole fraction of a . Sketch it.
- Derive the steady state profile of the temperature. Sketch it.
- If $x_{\alpha,L} > x_{\alpha,0}$, what is the relationship between $T(z=L)$ and $T(z=0)$?

On the next page, you will find a list of equations, some of which you will certainly find useful in solving this problem.

Note: This problem is much easier to solve if you assume that the heat capacities of α and β are equal. It is acceptable to make this assumption for this problem.

molar flux: $J_{\alpha}^* = -cD_{\alpha\beta} \frac{\partial x_{\alpha}}{\partial z}$	$J_{\beta}^* = -cD_{\alpha\beta} \frac{\partial x_{\beta}}{\partial z}$
for binary systems: $x_{\alpha} + x_{\beta} = 1$ $N = N_{\alpha} + N_{\beta}$ $\left(\frac{\partial x_{\alpha}}{\partial N_{\alpha}}\right)_{T,p,N_{\beta}} = \frac{1}{N_{\alpha} + N_{\beta}} - \frac{N_{\alpha}}{(N_{\alpha} + N_{\beta})^2} = \frac{N_{\beta}}{(N_{\alpha} + N_{\beta})^2} = x_{\beta} \frac{1}{N}$	$\frac{\partial x_{\beta}}{\partial z} = -\frac{\partial x_{\alpha}}{\partial z}$ $x_{\alpha} = \frac{N_{\alpha}}{N} = \frac{N_{\alpha}}{N_{\alpha} + N_{\beta}}$ $\left(\frac{\partial N}{\partial N_{\alpha}}\right)_{T,p,N_{\beta}} = 1$
$\left(\frac{\partial x_{\beta}}{\partial x_{\alpha}}\right)_{T,p,N_{\beta}} = \left(\frac{\partial \frac{N_{\beta}}{N_{\alpha} + N_{\beta}}}{\partial N_{\alpha}}\right)_{T,p,N_{\beta}} \quad \left(\frac{\partial N_{\alpha}}{\partial x_{\alpha}}\right)_{T,p,N_{\beta}} = -\frac{N_{\beta}}{(N_{\alpha} + N_{\beta})^2} \frac{1}{x_{\beta} \frac{1}{N}} = -1$	
pure component enthalpy $H_{\alpha}(T) = \int_{T_{\text{ref}}}^T C_{p,\alpha}(T) dT + H_{\alpha}(T_{\text{ref}})$	mixture enthalpy $H_{\text{mix}}(T, \underline{x}) = x_{\alpha} H_{\alpha}(T) + x_{\beta} H_{\beta}(T)$
partial molar enthalpy for an ideal mixture (derived) $\bar{H}_{\alpha}(T, \underline{x}) \equiv \left(\frac{\partial N H_{\text{mix}}}{\partial N_{\alpha}}\right)_{T,p,N_{\beta}} = N \left(\frac{\partial H_{\text{mix}}}{\partial N_{\alpha}}\right)_{T,p,N_{\beta}} + H_{\text{mix}} \left(\frac{\partial N}{\partial N_{\alpha}}\right)_{T,p,N_{\beta}}$ $\bar{H}_{\alpha}(T, \underline{x}) = N \left(\frac{\partial H_{\text{mix}}}{\partial x_{\alpha}}\right)_{T,p,N_{\beta}} \left(\frac{\partial x_{\alpha}}{\partial N_{\alpha}}\right)_{T,p,N_{\beta}} + H_{\text{mix}} = x_{\beta} \left(\frac{\partial H_{\text{mix}}}{\partial x_{\alpha}}\right)_{T,p,N_{\beta}} + H_{\text{mix}}$ $\left(\frac{\partial H_{\text{mix}}}{\partial x_{\alpha}}\right)_{T,p,N_{\beta}} = H_{\alpha}(T) + H_{\beta}(T) \left(\frac{\partial x_{\beta}}{\partial x_{\alpha}}\right)_{T,p,N_{\beta}} = H_{\alpha}(T) - H_{\beta}(T)$ $\bar{H}_{\alpha}(T, \underline{x}) = x_{\beta} [H_{\alpha}(T) - H_{\beta}(T)] + H_{\text{mix}} = (1 - x_{\alpha}) [H_{\alpha}(T) - H_{\beta}(T)] + H_{\text{mix}} = H_{\alpha}(T)$ $\bar{H}_{\beta}(T, \underline{x}) = x_{\alpha} [H_{\beta}(T) - H_{\alpha}(T)] + H_{\text{mix}} = \bar{H}_{\alpha}(T, \underline{x}) + [H_{\beta}(T) - H_{\alpha}(T)] = H_{\beta}(T)$	